

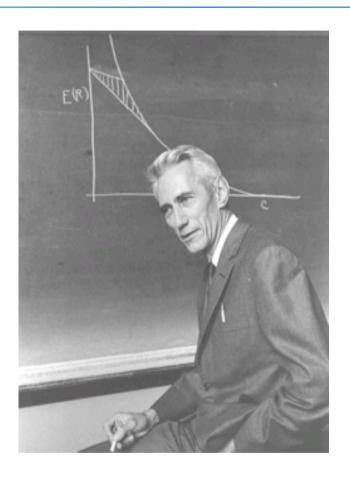
Information Rates for Phase Noise Channels (including fiber optic channels)

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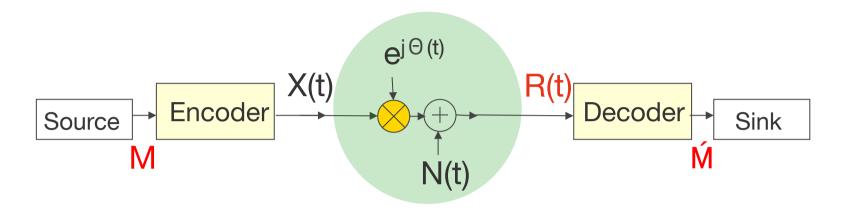




Claude Elwood Shannon Apr 30, 1916 – Feb 24, 2001

1) Phase Noise Models

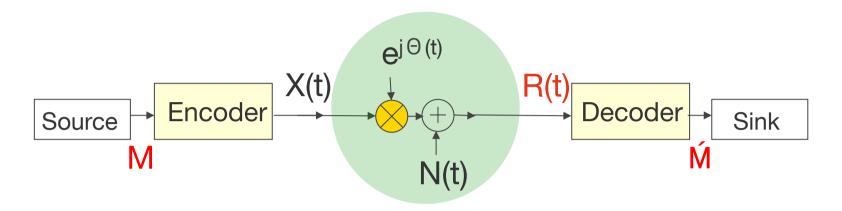




- Phase noise due to (1) oscillator instability; (2) fiber non-linearities
- Phase noise statistics:
 - phase-locked loops (PLLs) residual noise: von Mises/Tikhonov distribution
 - satellite (DVB-S2): white Gaussian process filtered by IIR filters
 - fiber-optic lasers: Wiener process
 - Raman amplification: large bandwidth Gaussian process

White Phase Noise





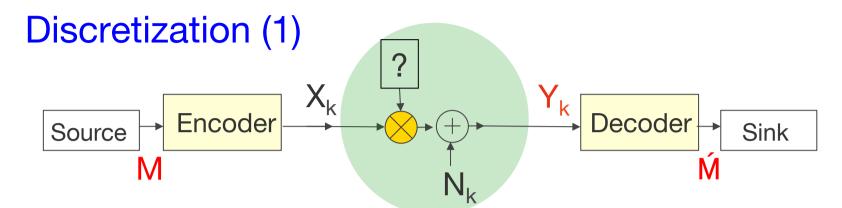
Simplified model (Barletta-Kramer, 2014)

$$R(t) = X(t) \cdot e^{j\Theta(t)} + N(t)$$

Θ(t) is white* and N(t) is white Gaussian* (both are idealizations)

- Motivation: phase noise bandwidth much larger than receiver bandwidth
- Mathematically: let $\{\emptyset_m(t)\}$ be an orthonormal basis of L²[0,T] and project X(t), N(t), and R(t) onto the $\emptyset_m(t)$





X(t) and N(t):

$$X(t) = \sum_{m=1}^{M} X_m \phi_m(t) \qquad N(t) = \sum_{m=1}^{\infty} N_m \phi_m(t)$$

Receiver:

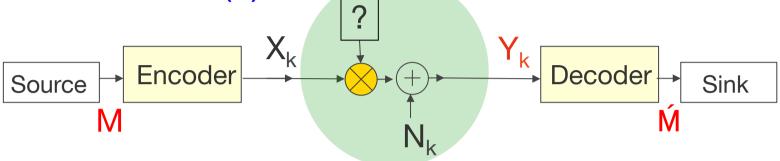
$$Y_{k} = \left\langle X(t)e^{j\Theta(t)} + N(t), \phi_{k}(t) \right\rangle$$

$$= \sum_{m=1}^{M} X_{m} \left\langle \phi_{m}(t)e^{j\Theta(t)}, \phi_{k}(t) \right\rangle + N_{k}$$

$$\Phi_{m,k}$$







Samples:

$$\Phi_{m,k} = \int_0^T \phi_m(t) e^{j\Theta(t)} \phi_k(t)^* dt$$

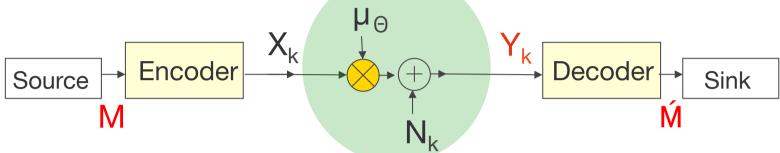
Barletta-Kramer, 2014:
Almost sure convergence for white phase noise with uncorrelated samples of process {e^{j © (t)}}

$$= \lim_{L \to \infty} \frac{1}{L} \sum_{i=1}^{L} \phi_m \left(t_i^{(L)} \right) e^{j\Theta\left(t_i^{(L)}\right)} \phi_k \left(t_i^{(L)} \right)^*, \quad t_i^{(L)} = \frac{(i-1)T}{L}$$

$$! \begin{cases} E\left[e^{j\Theta(t)}\right], & m = k \\ 0, & \text{else} \end{cases}$$







• Model*:
$$Y_k = X_k \cdot E[e^{j\Theta(t)}] + N_k$$

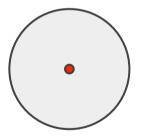
An AWGN channel (!) but with SNR penalty $|\mu_{\Theta}|^2$

- Penalty called spectral loss**: "lost" power is spread across all frequencies as white noise
- Proof: use Borel-Cantelli lemma with a classic trick and a simplified step via assumed boundedness of $\int |\phi_m(t)| \phi_k(t)^* dt$
- I expect this insight to be useful for fiber channels

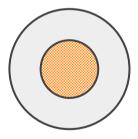


2) Fiber Channel(s)

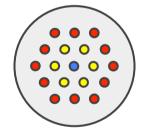
- Single-Mode Fiber (SMF): a small core that carries one mode of light
- Here one mode has 2 complex dimensions: two polarizations
- Theory papers often consider one complex dimension;
 the general case is interesting too of course (see below)
- In fact, a hot topic in the fiber community is MIMO fiber



Single-mode fiber (SMF)



Multi-mode fiber (MMF)

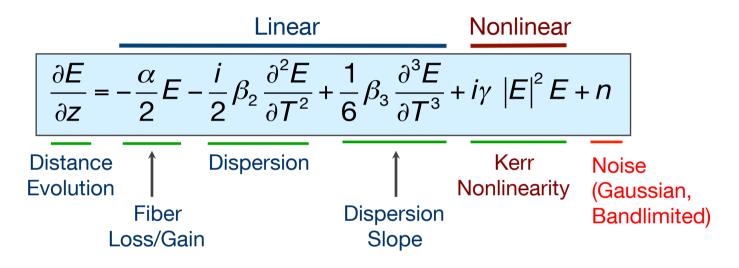


Multi-core fiber (MCF)

SMF Pulse Propagation Equation



 Maxwell's equations and low-order approximations* result in a generalized nonlinear Schrödinger equation (GNSE):



E: Electromagnetic field, function of z and T

z: Distance

T: Retarded time $t-\beta_1 z$

α: Fiber loss coefficient (~ 3 dB/15 km)

 β_1 : Inverse of group velocity

 β_2 : Fiber group velocity dispersion

 β_3 : Fiber dispersion slope (include if β_2 small)

 γ : Fiber nonlinear parameter ($n_2 \omega$)/(c A_{eff})

*n*₂: Fiber nonlinear coefficient

 ω : Angular frequency

c: Speed of light

A_{eff}: Fiber effective area

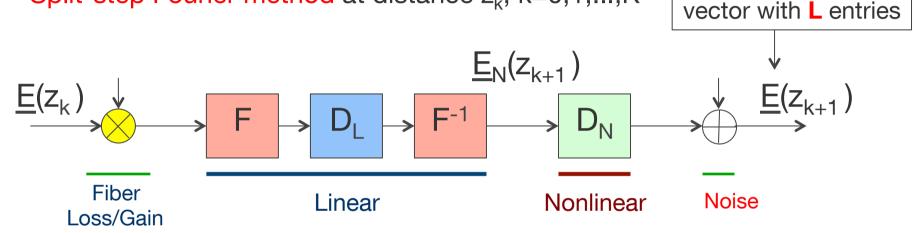
Figure courtesy of R.-J. Essiambre



time signal:

To simulate, split the fiber length z* into K small steps (Δz) and the time T into L small steps (Δt)

• Split-step Fourier method at distance z_k, k=0,1,...,K



- Ideal Raman amplification: removes the loss and adds noise
- F = Fourier transform
- D_L = diagonal matrix with fixed entries of unit amplitude (all-pass filter)
- $D_N =$ diagonal matrix with unit amplitude entries; the (ℓ, ℓ) -entry phase shift is proportional to the magnitude-squared of the ℓ^{th} entry of $\underline{E}_N(z_{k+1})$

3) Upper Bound



But First More IT Preliminaries

■ Consider a complex column vector $\underline{X} = \underline{X}_c + j \underline{X}_s$ with covariance and pseudo-covariance matrices

$$\mathbf{Q}_{\underline{X}} = E\left[\left(\underline{X} - E\left[\underline{X}\right]\right)\left(\underline{X} - E\left[\underline{X}\right]\right)^{\dagger}\right]$$

$$\tilde{\mathbf{Q}}_{\underline{X}} = E\left[\left(\underline{X} - E\left[\underline{X}\right]\right)\left(\underline{X} - E\left[\underline{X}\right]\right)^{T}\right]$$

- For interest: X is called proper if its pseudo-covariance matrix is 0
- Example: Consider a complex, zero-mean, scalar X = X_c + j X_s.
 X is proper if E[X_c²]=E[X_s²] and E[X_cX_s]=0.

Note: circularly symmetric X are proper, but proper X are not necessarily circularly symmetric (e.g. QAM signal sets)



Maximum Entropy

■ Maximum Entropy: consider the correlation matrix $\mathbf{R}_{\underline{X}} = \mathbb{E}[\underline{X} \ \underline{X}^{\dagger}]$ where \underline{X} has L entries. Then

$$h(\underline{X}) \leq \log[(\pi e)^L \det \mathbf{R}_{\underline{X}}]$$

with equality if and only if \underline{X} is Gaussian and proper (or circularly symmetric)

For a complex square matrix M we have

$$h(\mathbf{M} \underline{X}) = h(\underline{X}) + 2\log|\det(\mathbf{M})|$$

In particular, if **M** is unitary then $h(\mathbf{M} \times \mathbf{X}) = h(\mathbf{X})$



Entropy Power Inequality

Entropy Power:

$$V(\underline{X}) = e^{h(\underline{X})/L} / (\pi e)$$

Entropy Power Inequality: for independent X and Y we have

$$V(\underline{X} + \underline{Y}) \ge V(\underline{X}) + V(\underline{Y})$$

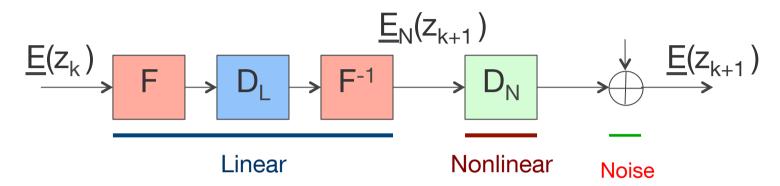
• Conditional version: for conditionally independent \underline{X} and \underline{Y} we have

$$V\left(\underline{X}|\underline{U}\right) = e^{h(\underline{X}|\underline{U})/L} / (\pi e)$$

$$V\left(\underline{X} + \underline{Y}|\underline{U}\right) \ge V\left(\underline{X}|\underline{U}\right) + V\left(\underline{Y}|\underline{U}\right)$$



Energy and Entropy Conservation



Main Observations

- The linear step conserves energy and entropy
- The non-linear step also conserves energy and entropy (the key result)

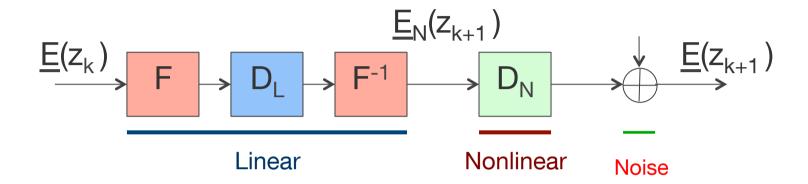
$$h\left(|a|e^{\int arg(a) + \int f(|a|)}\right) = h\left(|a|, arg(a) + f(|a|)\right) + E\left[\log|a|\right]$$

$$= h\left(|a|\right) + h\left(arg(a) + f(|a|) | |a|\right) + E\left[\log|a|\right] = h(a)$$

$$h\left(|a|, arg(a)\right)$$



Energy Recursion

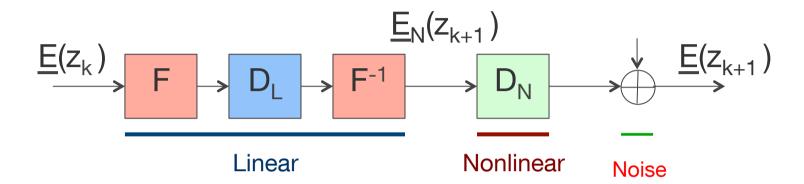


• Energy after K steps: Energy_{Launch} + KN . We thus have:

$$h(\underline{E}(z_{\kappa})) \leq \log[(\pi e)^{L} \det(\mathbf{R}(\underline{E}(z_{\kappa})))]$$
 ... maximum entropy
$$\leq \sum_{i=1}^{L} \log[\pi e \, R_{i,i}(\underline{E}(z_{\kappa}))]$$
 ... Hadamard's inequality
$$\leq L \cdot \log[\pi e \, (Energy_{Launch} + KN)/L]$$
 ... Jensen's inequality



Entropy Recursion



Entropy recursion:

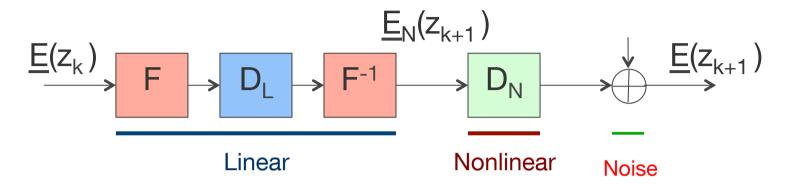
$$V\left(\underline{E}(Z_{k+1})|\underline{E}(Z_0)\right) \ge V\left(\underline{E}(Z_k)|\underline{E}(Z_0)\right) + N/L$$

We thus have:

$$V(\underline{E}(z_K)|\underline{E}(z_0)) \ge KN/L$$

or $h(\underline{E}(z_K)|\underline{E}(z_0)) \ge L\log(\pi e \, KN/L)$





So for every step we have:

- Signal energy grows by the noise variance: can upper bound h(<u>E(z_k)</u>)
- Entropy power grows by at least the noise variance: can lower bound h($\underline{E}(z_K) \mid \underline{E}(z_0)$)
- Result*:

$$I(\underline{E}(z_0);\underline{E}(z_K)) = h(\underline{E}(z_K)) - h(\underline{E}(z_K)|\underline{E}(z_0))$$

$$\leq L \cdot \log(1 + SNR)$$



$$\Rightarrow \frac{1}{L}I(\underline{E}(z_0);\underline{E}(z_{\kappa})) \leq \log(1+SNR)$$

- Let B = $1/\Delta t$ be the "bandwidth" of the simulation
- So L = $T/\Delta t$ = TB is the time-bandwidth product
- The spectral efficiency is thus bounded by

$$\eta \le \log(1 + SNR)$$
 [bits/sec/Hz]



5) Conclusions

- Spectral efficiency of (an idealized model of) SMF with linear polarization is ≤ log(1+SNR)
- 2) Many extensions are possible:
 - lumped amplification, 3rd-order dispersion, delayed Kerr effect
 - uniform loss, linear filters (for capacity results)
 - MIMO fiber (MMF or MCF)
- 3) More difficult:
 - better bounds and understanding at high SNR
 - frequency-dependent loss, dispersion, non-linearity